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YANG MODEL REVISITED¹

A long time ago, C.N. Yang proposed a generalization of the Snyder model to the case of a curved background spacetime, based on an algebra isomorphic to $o(1, 5)$ which includes, as sub-algebras both the Snyder and the de Sitter algebras. His proposal can, therefore, be interpreted as a model of noncommutative curved spacetime, and could be useful for relating physics on very small and very large scales. We review this model and some recent progress concerning its generalizations and its interpretation in the framework of Hopf algebras. We also report some possibilities to relate it to more phenomenological aspects.

Keywords: noncommutative geometry, de Sitter spacetime, Yang model.

1. Quantum Gravity and Noncommutative Geometry

At present, no complete theory of quantum gravity is available. However, it is known that the predictions of quantum mechanics and general relativity imply the existence of a minimal measurable length of the size of the Planck length $L_P = \sqrt{\frac{\hbar G}{c^3}} \sim 10^{-33}$ cm [1, 2]. Hence, it is natural to believe that the properties of spacetime on this scale must be rather different from those we observe in everyday experience.

Among the proposals for a model of spacetime on the Planck scale, noncommutative geometry plays a relevant role [3]. Noncommutative geometry is based on the assumption that the components of the position operator do not commute, leading to the impossibility of localizing a particle exactly. Among various approaches to this field, a relevant one is the Hopf algebra formalism [4–6], which is suitable for the description of the symmetries of the geometry.

Noncommutative geometries are usually defined on a flat spacetime, but their extension to curved spacetimes has earned some interest recently because of possible implications for astrophysical observations, like time delay of photons from distant sources [7]. However, also the formal aspects of this extension are noticeable, in particular, the relations be-

tween the curvature of spacetime and of momentum space. Moreover, these models relate the properties of spacetime in microscopic and macroscopic scales.

The first model of this kind was proposed by C.N. Yang already in 1947 [8]. In the present paper, we review this framework and discuss some recent progresses and generalizations.

2. The Snyder Algebra

We start by summarizing the main features of the Snyder model, from which Yang took inspiration. In 1947 Snyder proposed the first model of noncommutative geometry [9]. His aim was to construct a theory that included a fundamental length without breaking the Lorentz invariance. This purpose was realized by deforming the commutation relations of the Heisenberg algebra.

In fact, the model was defined through an algebra that, besides the deformed Heisenberg algebra generated by positions \hat{x}_μ and momenta \hat{p}_μ , contained the Lorentz algebra with generators $J_{\mu\nu}$,

$$\begin{aligned} [\hat{x}_\mu, \hat{x}_\nu] &= i\beta J_{\mu\nu}, & [\hat{p}_\mu, \hat{p}_\nu] &= 0, \\ [\hat{x}_\mu, \hat{p}_\mu] &= i(\eta_{\mu\nu} + \beta\hat{p}_\mu\hat{p}_\nu), \\ [J_{\mu\nu}, J_{\rho\sigma}] &= i(\eta_{\mu\rho}J_{\nu\sigma} - \eta_{\mu\sigma}J_{\nu\rho} + \\ &+ \eta_{\nu\rho}J_{\mu\sigma} - \eta_{\nu\sigma}J_{\mu\rho}), \\ [J_{\mu\nu}, \hat{p}_\lambda] &= i(\eta_{\mu\lambda}\hat{p}_\nu - \eta_{\lambda\nu}\hat{p}_\mu), \\ [J_{\mu\nu}, \hat{x}_\lambda] &= i(\eta_{\mu\lambda}\hat{x}_\nu - \eta_{\nu\lambda}\hat{x}_\mu). \end{aligned} \quad (1)$$

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In particular, the components of the position operator \hat{x}_μ do not commute among themselves. The coupling constant β has dimension of inverse mass square and may be identified with $1/M_P^2 = L_P^2$, M_P being the Planck mass. Note that β can take both signs. If it is positive, the bound $p^2 < 1/\beta$ must hold, and, hence, a maximal mass $1/\sqrt{\beta}$ is predicted.

Snyder also showed that his algebra can be realized on canonical phase space of coordinates x and p , with Lorentz generators defined as $J_{\mu\nu} = x_\mu p_\nu - x_\nu p_\mu$.

In contrast with most common models of noncommutative geometry, the commutators are not linear in the phase space variables: this allows them to be compatible with a linear action of the Lorentz symmetry, so that the Poincaré algebra is not deformed. However, translations (generated by the \hat{p}_μ) act in a nonlinear way on position variables.

The Snyder model can be interpreted as describing the flat spacetime with a curved momentum space. In fact, the subalgebra generated by $J_{\mu\nu}$ and \hat{x}_μ is isomorphic to the de Sitter algebra $o(1, 4)$, and, hence, the Snyder momentum space has the same geometry as the de Sitter spacetime.

3. The Yang Algebra

Soon after the publication of Snyder's paper, Yang proposed a generalization where also the momentum variables do not commute, as is the case in the de Sitter spacetime [8].

The algebra is isomorphic to $o(1, 5)$, with 15 generators,

$$\begin{aligned} [\hat{x}_\mu, \hat{x}_\nu] &= i\beta J_{\mu\nu}, & [\hat{p}_\mu, \hat{p}_\nu] &= i\alpha J_{\mu\nu}, \\ [\hat{x}_\mu, \hat{p}_\nu] &= i\eta_{\mu\nu} K, \\ [J_{\mu\nu}, J_{\rho\sigma}] &= i(\eta_{\mu\rho} J_{\nu\sigma} - \eta_{\mu\sigma} J_{\nu\rho} + \\ &+ \eta_{\nu\rho} J_{\mu\sigma} - \eta_{\nu\sigma} J_{\mu\rho}), \\ [J_{\mu\nu}, \hat{p}_\lambda] &= i(\eta_{\mu\lambda} \hat{p}_\nu - \eta_{\lambda\nu} \hat{p}_\mu), \\ [J_{\mu\nu}, \hat{x}_\lambda] &= i(\eta_{\mu\lambda} \hat{x}_\nu - \eta_{\nu\lambda} \hat{x}_\mu), & [J_{\mu\nu}, K] &= 0, \\ [K, \hat{x}_\mu] &= i\beta \hat{p}_\mu, & [K, \hat{p}_\mu] &= -i\alpha \hat{x}_\mu. \end{aligned} \tag{2}$$

The deformation parameter α has dimension of inverse length square and may be identified with the cosmological constant, while β is the same as in the Snyder model. Notice that both α and β can take positive or negative values, giving rise to models with very different physical properties [10]. For example, positive α enforces the bound $\hat{x}^2 < 1/\alpha^2$ and analogously, for positive β , $\hat{p}^2 < 1/\beta^2$.

The Yang algebra contains as subalgebras both the de Sitter and the Snyder algebras, and, therefore, describes a noncommutative model in a spacetime of constant curvature. In order to close the algebra, Yang had to introduce a new generator K which rotates positions into momenta, but whose physical interpretation is not evident.

The algebra (2) is invariant under a generalized Born duality [11, 12],

$$\begin{aligned} \alpha &\leftrightarrow \beta, & \hat{x}_\mu &\rightarrow -\hat{p}_\mu, & \hat{p}_\mu &\rightarrow \hat{x}_\mu, \\ J_{\mu\nu} &\leftrightarrow J_{\mu\nu}, & K &\leftrightarrow K. \end{aligned} \tag{3}$$

The isomorphism with the $o(1, 5)$ algebra can be obtained by identifying

$$M_{\mu\nu} = J_{\mu\nu}, \quad M_{\mu 4} = \hat{x}_\mu, \quad M_{\mu 5} = \hat{p}_\mu, \quad M_{45} = K, \tag{4}$$

where M_{AB} ($A, B = 0, \dots, 5$) are the generators of $o(1, 5)$.

There also exists a different generalization of the Snyder algebra onto a curved space, known as triply special relativity (TSR), that does not include K , but is nonlinear [13].

In particular, in that case, the deformed Heisenberg subalgebra takes the form

$$\begin{aligned} [\hat{x}_\mu, \hat{x}_\nu] &= i\beta J_{\mu\nu}, & [\hat{p}_\mu, \hat{p}_\nu] &= i\alpha J_{\mu\nu}, \\ [\hat{x}_\mu, \hat{p}_\nu] &= i\left(\eta_{\mu\nu} + \alpha \hat{x}_\mu \hat{x}_\nu + \beta \hat{p}_\mu \hat{p}_\nu + \right. \\ &+ \left. \sqrt{\alpha\beta}(\hat{x}_\mu \hat{p}_\nu + \hat{p}_\mu \hat{x}_\nu - J_{\mu\nu})\right), \end{aligned} \tag{5}$$

and one can interpret the phase space as a coset space $SO(1, 5)/SO(1, 3) \times O(2)$.

Triply special relativity theory has been described as a deformation of the Galilei group in [14, 15], and its Hamiltonian formulation has been investigated in [16, 17]. A framework unifying it with the Yang model has been proposed in [18].

Two interpretations of the Yang algebra are possible:

- Take the algebra as it is, with its 15 generators [19, 20]. This allows one to construct the Hopf algebra structure, with the related star product, twist, *etc.* exploiting the results known for generic orthogonal algebras [21]. However, in this case, one has to consider an extended phase space with scalar and tensorial degrees of freedom, whose interpretation is not obvious.

- Take a nonlinear realization on canonical phase space spanned by x_μ and p_μ , with $J_{\mu\nu} = x_\mu p_\nu - x_\nu p_\mu$ and $K = K(x, p)$ [22, 23], in analogy with the representation given by Snyder for his model [9]. In this case, the interpretation is easier, and one can include the Yang model in the same family of nonlinear realizations of $o(1, 5)$ as TSR, identifying the phase space with a coset space. However, one can no longer define a Hopf algebra, star products, etc.

4. Realizations of the Yang–Poisson Model

We start following the second route, and discussing the classical limit of the Yang model, in which commutators are replaced by Poisson brackets. This makes the investigation much easier because of the absence of ordering problems. We call this limit Yang–Poisson model.

We have [24]

$$\begin{aligned} \{\hat{x}_\mu, \hat{x}_\nu\} &= \beta J_{\mu\nu}, & \{\hat{p}_\mu, \hat{p}_\nu\} &= \alpha J_{\mu\nu}, \\ \{\hat{x}_\mu, \hat{p}_\nu\} &= \eta_{\mu\nu} K, & & \\ \{K, \hat{x}_\mu\} &= \beta \hat{p}_\mu, & \{K, \hat{p}_\mu\} &= -\alpha \hat{x}_\mu, \end{aligned} \tag{6}$$

and look for an expression of $K(x, p)$ that satisfies the previous Poisson brackets

$$\begin{aligned} \hat{x}_\mu &= f(p^2, z) x_\mu, & \hat{p}_\mu &= g(x^2, z) p_\mu, \\ K &= K(x^2, p^2, z), \end{aligned} \tag{7}$$

where $z = x \cdot p$, and f and g are functions to be determined.

The only nontrivial brackets to be checked are those of the deformed Heisenberg algebra, which give rise to partial differential equations. The equations derived from the x - x and p - p brackets have solutions

$$f = \sqrt{1 - \beta p^2 + \phi_1(z)}, \quad g = \sqrt{1 - \alpha x^2 + \phi_2(z)}, \tag{8}$$

with arbitrary functions ϕ_1 and ϕ_2 , while the x - p brackets give

$$\phi_1 \phi_2 + \phi_1 + \phi_2 = \alpha \beta z^2, \quad K = fg, \tag{9}$$

with solution depending on one parameter c ,

$$\begin{aligned} \phi_1(z) &= \frac{\sqrt{1 + 4c(1-c)z^2} - 1}{2(1-c)}, \\ \phi_2(z) &= \frac{\sqrt{1 + 4c(1-c)z^2} - 1}{2c}. \end{aligned} \tag{10}$$

Then,

$$\begin{aligned} \hat{x}_\mu &= \sqrt{1 - \beta p^2 + \phi_1(z)} x_\mu, \\ \hat{p}_\mu &= \sqrt{1 - \alpha x^2 + \phi_2(z)} p_\mu, \end{aligned} \tag{11}$$

and

$$K = \sqrt{[1 - \beta p^2 + \phi_1(z)][1 - \alpha x^2 + \phi_2(z)]}. \tag{12}$$

A particularly interesting solution is obtained by assuming symmetry under the exchange of x and p , as is natural in view of the Born duality of the model. In this case, $\phi_1 = \phi_2 = \phi$, i.e., $c = \frac{1}{2}$, and we obtain

$$\phi = \sqrt{1 + \alpha \beta z^2} - 1, \tag{13}$$

and then

$$\begin{aligned} \hat{x}_\mu &= \sqrt{\sqrt{1 + \alpha \beta z^2} - \beta p^2} x_\mu, \\ \hat{p}_\mu &= \sqrt{\sqrt{1 + \alpha \beta z^2} - \alpha x^2} p_\mu. \end{aligned} \tag{14}$$

This gives an exact realization of the Yang–Poisson model, symmetric for $x \leftrightarrow p$ and $\alpha \leftrightarrow \beta$. One can also write K in terms of the original variables, as

$$\begin{aligned} K &= \\ &= \sqrt{\frac{1 - \alpha \hat{x}^2 - \beta \hat{p}^2}{2} \left(1 + \sqrt{1 - 4\alpha\beta \frac{\hat{x}^2 \hat{p}^2 - (\hat{x} \cdot \hat{p})^2}{(1 - \alpha \hat{x}^2 - \beta \hat{p}^2)^2}} \right)}. \end{aligned} \tag{15}$$

More general realizations can also be found [24, 25].

5. Realizations of the Quantum Yang Model

In the quantum case, finding a realization is more difficult, and no general closed form is known. One has, therefore, to resort to a perturbative calculation in the coupling parameters α and β .

A perturbative calculation up to fourth order was first performed in [22, 23]. More elaborated methods were used in [20], where also representations in extended phase space were investigated.

Realizations can be found order by order, making an ansatz that includes the most general Lorentz-covariant terms. For example, in first order in α and β , one can find an Hermitian realization setting

$$\begin{aligned} \hat{x}_\mu &= x_\mu + (a_1 \sqrt{\alpha \beta} x_\mu x \cdot p + a_2 \beta x_\mu p^2 + \\ &+ a_3 \beta p_\mu p \cdot x + a_4 \sqrt{\alpha \beta} p_\mu x^2 + \text{h.c.}), \\ \hat{p}_\mu &= p_\mu + (b_1 \sqrt{\alpha \beta} p_\mu p \cdot x + b_2 \alpha p_\mu x^2 + \\ &+ b_3 \alpha x_\mu x \cdot p + b_4 \sqrt{\alpha \beta} x_\mu p^2 + \text{h.c.}), \\ K &= 1 + (h_1 \alpha x^2 + h_2 \sqrt{\alpha \beta} x \cdot p + h_3 \beta p^2 + \text{h.c.}), \end{aligned} \tag{16}$$

and substituting in (1), obtaining constraints between the free parameters.

Remarkably, the simplest solution of these constraints symmetric in \hat{x} and \hat{p} , is given in the second order by [22, 23]

$$\begin{aligned} \hat{x}_\mu &= x_\mu - \frac{\beta}{4}x_\mu p^2 - \frac{\beta^2}{16}x_\mu p^4 + \\ &+ \frac{\alpha\beta}{8}x_\mu x \cdot p p \cdot x + \text{h.c.}, \\ \hat{p}_\mu &= p_\mu - \frac{\alpha}{4}p_\mu x^2 - \frac{\alpha^2}{16}p_\mu x^4 + \\ &+ \frac{\alpha\beta}{8}p_\mu p \cdot x x \cdot p + \text{h.c.}, \end{aligned} \tag{17}$$

with

$$K = 1 - \frac{1}{2}(\alpha x^2 + \beta p^2) - \frac{1}{8}(\alpha x^2 - \beta p^2)^2 + \frac{\alpha\beta}{2}x \cdot p p \cdot x. \tag{18}$$

Of course, this is nothing but the expansion of the Yang–Poisson result (14) in powers of α and β . However, going to higher orders in α and β , one finds corrections with respect to the classical solution [20].

6. Star Product for the Yang Algebra

One of the most useful frameworks for the investigation of noncommutative geometry is that of Hopf algebras [4–6]. This formalism implies the definition of a coalgebraic structure including a coproduct and an antipode.

Such formalism can be applied also to the Yang model, provided one takes all its generators M_{AB} as primary variables [20, 26], and, therefore, use a realization on the extended phase space [25]. One can then use the isomorphism (4) of the Yang algebra with $o(1, 5)$ and the general results of [21], where the Hopf algebra related to general orthogonal groups was computed. Notice that, to this end, it is necessary to define “momenta” s^{AB} canonically conjugated to the variables M_{AB} , whose physical interpretation is not obvious.

We shall not illustrate the details here, referring the reader to the original literature. We only recall that, due to the noncommutativity, the addition law of the momenta s^{AB} is deformed [4–6]. Using the Hopf algebra formalism, this deformation can be expressed by means of a star product, whose knowledge can be useful, for example, in the construction of a quantum field theory.

In our case, the star product for plane waves can be defined as

$$e^{\frac{i}{2}s^{AB}M_{AB}} \star e^{\frac{i}{2}t^{CD}M_{CD}} = e^{\frac{i}{2}\mathcal{D}^{AB}(s,t)M_{AB}}, \tag{19}$$

where s_{AB} and t_{AB} are antisymmetric tensors that describe the “momenta” conjugated to the primary variables M_{AB} , and \mathcal{D}^{AB} encodes the deformed addition law.

Using a Weyl realization of the algebra, the star product can be calculated perturbatively [26].

It may be useful to explicitly write down the four-dimensional expression of $\mathcal{D}^{AB}(s, t)$ in the first order: setting $\mathcal{D}^\mu = \mathcal{D}^{\mu 4}$, $\bar{\mathcal{D}}^\mu = \mathcal{D}^{\mu 5}$, $\mathcal{D} = \mathcal{D}^{45}$, one has

$$\begin{aligned} \mathcal{D}^{\mu\nu}(s, t) &= s^{\mu\nu} + t^{\mu\nu} - \frac{1}{2}(s^{\mu\lambda}t^\nu{}_\lambda + \beta s^{\mu t^\nu} + \\ &+ \alpha \bar{s}^\mu \bar{t}^\nu + \sqrt{\alpha\beta}(s^\mu \bar{t}^\nu + \bar{s}^\mu t^\nu) - (\mu \leftrightarrow \nu)), \\ \mathcal{D}^\mu(s, t) &= s^\mu + t^\mu - \frac{1}{2}(s^{\mu\lambda}t_\lambda - t^{\mu\lambda}s_\lambda + \\ &+ \sqrt{\alpha\beta}(s^{\mu t} - st^\mu) + \alpha(\bar{s}^\mu t - s\bar{t}^\mu)), \\ \bar{\mathcal{D}}^\mu(s, t) &= \bar{s}^\mu + \bar{t}^\mu - \frac{1}{2}(s^{\mu\lambda}\bar{t}_\lambda - \bar{s}_\lambda t^{\mu\lambda} - \\ &- \sqrt{\alpha\beta}(\bar{s}^\mu t - s\bar{t}^\mu) + \beta(s^{\mu t} - st^\mu)), \\ \mathcal{D}(s, t) &= s + t - \frac{1}{2}(s^\lambda \bar{t}_\lambda - \bar{s}^\lambda t_\lambda), \end{aligned} \tag{20}$$

where $s^{\mu\nu}$, $s^\mu = s^{\mu 4}$, $\bar{s}^\mu = s^{\mu 5}$, $s = s^{45}$ are the four-dimensional components of s^{AB} , conjugated to $J_{\mu\nu}$, x_μ , p_μ and K , respectively. It is clear from these expressions that the four-dimensional components of the momenta mix in the star product.

7. Generalizations

The Yang algebra can be generalized in several ways. The simplest one is to admit negative values of the deformation parameters α or β , obtaining models based on the algebras $o(2, 4)$ or $o(3, 3)$ [25]. These models have different physical properties from those with positive α and β [10].

Another simple generalization was given in [27], where linear combinations of the generators \hat{x} and \hat{p} by a parameter γ were introduced. In [18, 22, 23] still more general definitions of the Yang model in extended phase space were proposed, unifying it with TSR. In addition supersymmetric extension of the Yang algebra was investigated in [28].

Finally, the possibility to include κ -deformations of both position and momentum space with parameters

a_μ and b_μ into the model of Ref. [27] was studied in [26, 29]. The generalized algebra reads

$$\begin{aligned}
[x_\mu, x_\nu] &= i(\beta J_{\mu\nu} + a_\mu x_\nu - a_\nu x_\mu), \\
[p_\mu, p_\nu] &= i(\alpha J_{\mu\nu} + b_\mu p_\nu - b_\nu p_\mu), \\
[x_\mu, p_\nu] &= i(\eta_{\mu\nu}\hbar + b_\mu x_\nu - a_\nu p_\mu + \gamma J_{\mu\nu}), \\
[J_{\mu\nu}, x_\lambda] &= i(\eta_{\mu\lambda}x_\nu - \eta_{\nu\lambda}x_\mu + a_\mu J_{\lambda\nu} - a_\nu J_{\lambda\mu}), \\
[J_{\mu\nu}, p_\lambda] &= i(\eta_{\mu\lambda}p_\nu - \eta_{\nu\lambda}p_\mu + b_\mu J_{\lambda\nu} - b_\nu J_{\lambda\mu}), \\
[J_{\mu\nu}, K] &= i(b_\nu x_\mu - b_\mu x_\nu - a_\nu p_\mu + a_\mu p_\nu), \\
[K, x_\mu] &= i(\beta p_\mu - \gamma x_\mu - a_\mu K), \\
[K, p_\nu] &= i(-\alpha x_\mu + \gamma p_\mu + b_\mu K).
\end{aligned} \tag{21}$$

This algebra is still isomorphic to $o(1,5)$, but now, the action of the Lorentz group is deformed. It is, however, still possible to obtain the coalgebra associated to it [26].

8. Applications

A physical consequence of the Yang model is a deformation of the Heisenberg uncertainty relations [10]. In fact, in its nonrelativistic 3-dimensional limit,

$$\Delta x_i \Delta p_j \geq \frac{1}{2} |\langle [x_i, p_j] \rangle| = \frac{1}{2} |\langle K \rangle| \delta_{ij}. \tag{22}$$

Clearly, in a phase space realization with $K = K(x, p)$ this gives rise to a generalized uncertainty principle, similar to the one valid for TSR [30].

For example, let us consider a one-dimensional toy model: at leading order in \hbar , α and β , one can use the realization (14), obtaining¹

$$\Delta \hat{x} \Delta \hat{p} \geq \frac{\hbar}{2} \sqrt{1 - \alpha(\Delta \hat{x})^2 - \beta(\Delta \hat{p})^2}. \tag{23}$$

It follows that

$$\begin{aligned}
(\Delta \hat{x}) &\geq \hbar \sqrt{\frac{1 - \beta(\Delta \hat{p})^2}{\hbar^2 \alpha + 4(\Delta \hat{p})^2}}, \\
(\Delta \hat{p}) &\geq \hbar \sqrt{\frac{1 - \alpha(\Delta \hat{x})^2}{\hbar^2 \beta + 4(\Delta \hat{x})^2}}.
\end{aligned} \tag{24}$$

Clearly, the uncertainty relations will strongly depend on the sign of the coupling constants. For example, for $\alpha, \beta > 0$, the uncertainties satisfy the bounds $0 \leq \Delta p \leq 1/\sqrt{\beta}$, and $0 \leq \Delta x \leq 1/\sqrt{\alpha}$, as is obvious

¹ In this section, we reintroduce explicitly \hbar .

because of the limited range of variation of x and p in this case. For $\alpha, \beta < 0$, instead, one gets the interesting relations $\Delta p \geq \hbar\sqrt{|\alpha|}/2$, and $\Delta x \geq \hbar\sqrt{|\beta|}/2$. A more complicated behaviour occurs if α and β have different sign [10].

One may also calculate corrections to the dynamics of simple models due to the nontrivial symplectic structure, with possible applications to astrophysical observations, again in analogy with the results of [30] for TSR.

Finally, a more ambitious goal would be to build a quantum field theory based on this framework. This would presumably require the knowledge of the star product and hence the use of an extended phase space. In this case, the physical interpretation of the additional coordinates needs to be clarified.

9. Final Remarks

We have reviewed the main properties of a model of noncommutative geometry on a curved background proposed by C.N. Yang in the 1940s, and have reported some developments and generalizations obtained in the last years.

The physical relevance of these models is that they somehow relate the description of spacetime at extremely small distances and on cosmological scales, and that they may be connected to a low-energy limit of quantum gravity [13].

Some applications have been investigated, however, much remains to be understood concerning the physical interpretation and the phenomenological predictions of these models.

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1. L.J. Garay. Quantum gravity and minimum length. *Int. J. Mod. Phys. A* **10**, 145 (1995).
2. S. Hossenfelder. Minimal length scale scenarios for quantum gravity. *Liv. Rev. Rel.* **16**, 2 (2013).
3. S. Doplicher, K. Fredenhagen, J.E. Roberts. The quantum structure of spacetime at the Planck scale and quantum fields. *Commun. Math. Phys.* **172**, 187 (1995).
4. S. Majid. Algebraic approach to Quantum Gravity II: noncommutative spacetime. In: *Approaches to Quantum Gravity*. Edited by D. Oriti (Cambridge Univ. Press, 2009), p. 466.

5. J. Madore. *An Introduction to Noncommutative Geometry and Its Physical Applications* (Cambridge Univ. Press, 1995).
6. M. Arzano M., J. Kowalski-Glikman. *Deformation of Spacetime Symmetries – Gravity, Group-Valued Momenta, and Noncommutative Fields* (Springer-Verlag, 2021).
7. G. Rosati, G. Amelino-Camelia, A. Marciano, M. Matassa. Planck-scale-modified dispersion relations in FRW space-time. *Phys. Rev. D* **92**, 124042 (2015).
8. C.N. Yang. On Quantized space-time. *Phys. Rev.* **72**, 874 (1947).
9. H.S. Snyder. Quantized space-time. *Phys. Rev.* **71**, 38 (1947).
10. S. Meljanac, S. Mignemi. in preparation.
11. M. Born. Reciprocity theory of elementary particles. *Rev. Mod. Phys.* **21** 463 (1949).
12. H.G. Guo, C.G. Huang, H.T. Wu. Yang's model as triply special relativity and the Snyder's model-de Sitter special relativity duality. *Phys. Lett. B* **663** 270 (2008).
13. J. Kowalski-Glikman, L. Smolin. Triply special relativity. *Phys. Lett. D* **70**, 065020 (2004).
14. C. Chryssomakolos, E. Okon. Linear form of 3-scale special relativity algebra and the relevance of stability. *Int. J. Mod. Phys. D* **13**, 1817 (2004).
15. A. Das, O.C.W. Kong. Physics of quantum relativity through a linear realization. *Phys. Rev. D* **73**, 124029 (2006).
16. S. Mignemi. The Snyder model and quantum field theory. *Class. Quantum Grav.* **26**, 245020 (2009).
17. R. Banerjee, K. Kumar, D. Roychowdhury. Symmetries of Snyder–de Sitter space and relativistic particle dynamics. *J. High Energ. Phys.* **1103**, 060 (2011).
18. S. Meljanac, R. Štrajn. Deformed quantum phase spaces, realizations, star products and twists. *Symmetry, Integrability and Geometry: Methods and Applications* **18**, 022 (2022).
19. J. Lukierski, S. Meljanac, S. Mignemi, A. Pachol. Quantum perturbative solutions of extended Snyder and Yang models with spontaneous symmetry breaking. *Phys. Lett. B* **847**, 138261 (2023).
20. T. Martinić-Bilać, S. Meljanac, S. Mignemi. Hermitian realizations of the Yang model. *J. Math. Phys.* **64**, 122302 (2023).
21. S. Meljanac, T. Martinić-Bilać, S. Krešić-Jurić. Generalized Heisenberg algebra applied to realizations of the orthogonal, Lorentz, and Poincaré algebras and their dual extensions. *J. Math. Phys.* **61**, 051705 (2020).
22. S. Meljanac, S. Mignemi. Generalizations of Snyder model to curved spaces. *Phys. Lett. B* **833**, 137289 (2022).
23. S. Meljanac, S. Mignemi. Noncommutative Yang model and its generalizations. *J. Math. Phys.* **64**, 023505 (2023).
24. S. Meljanac, S. Mignemi. Realizations of the Yang–Poisson model on canonical phase space. *Int. J. Mod. Phys. A* **38**, 2350182 (2023).
25. T. Martinić-Bilać, S. Meljanac, S. Mignemi. Generalized Yang–Poisson models on canonical phase space. *Symmetry, Integrability and Geometry: Methods and Applications* **20**, 049 (2024).
26. T. Martinić-Bilać, S. Meljanac, S. Mignemi. Realizations and star-product of doubly κ -deformed Yang models. arXiv:2404.01792.
27. V.V. Khruschev, A.N. Leznov. The relativistic invariant Lie algebra for the kinematical observables in quantum space-time. *Grav. Cosmol.* **9**, 159 (2003).
28. J. Lukierski, M. Woronowicz. Spinorial Snyder and Yang models from superalgebras and noncommutative quantum superspaces. *Phys. Lett. B* **824**, 136783 (2021).
29. J. Lukierski, S. Meljanac, S. Mignemi, A. Pachol. From Snyder space-times to doubly κ -dependent Yang quantum phase spaces and their generalizations. *Phys. Lett. B* **854**, 138729 (2024).
30. S. Mignemi. Classical and quantum mechanics of the non-relativistic Snyder model in curved space. *Class. Quantum Grav.* **29**, 215019 (2012).

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С. Мінємі

ПЕРЕГЛЯД МОДЕЛІ ЯНГА

Дуже давно С.Н. Янг запропонував узагальнення моделі Снайдера на випадок викривленого фонового простору-часу на основі алгебри, ізоморфної $o(1,5)$, яка включає в себе, як підалгебри, і алгебру Снайдера, і алгебру де Сіттера. Тому його пропозицію можна сприймати як модель некомутативного викривленого простору-часу, і вона може бути корисною для того, щоб пов'язати фізику дуже малих і дуже великих масштабів. Ми розглядаємо цю модель і деякі останні досягнення, які стосуються її узагальнень та її інтерпретації в рамках алгебр Хопфа. Ми також обговорюємо деякі можливості пов'язати її з більш феноменологічними аспектами.

Ключові слова: некомутативна геометрія, простір-час де Сіттера, модель Янга.